# Electron-electron biwire systems





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Model

electron density n in each wire, coupling parameter  $r_s = 1/(2na_{\rm B})$ , Bohr radius  $a_{\rm B} = \epsilon/(e^2 m_e)$ , Hartree atomic units ( $\hbar = |e| = m_e = \epsilon = 1$ )





## Momentum distribution

fitting n(k) within the range  $|k - k_F| < \epsilon k_F$ 

$$n(k) = \frac{1}{2\pi} \left\langle \int \frac{\Psi_T(r)}{\Psi_T(x_1)} e^{ik(x_1 - r)} dr \right\rangle = n(k_F) + A[\operatorname{sign}(k - k_F)]|k - k_F|^{\alpha}$$

Luttinger parameter  $K_{\rho} = 1 + 2\alpha - 2\sqrt{\alpha + \alpha^2}$ 



$$\hat{H} = -\frac{1}{2} \sum_{i=1}^{N} \left( \frac{\partial^2}{\partial \mathbf{x}_{i,1}^2} + \frac{\partial^2}{\partial \mathbf{x}_{i,2}^2} \right) + \sum_{i < j} \left\{ V(x_{ij,1}) + V(x_{ij,2}) \right\} + \sum_{i,j} \frac{1}{\sqrt{|\mathbf{x}_{i,1} - \mathbf{x}_{j,2}|^2 + d^2}} + NV_{\text{Machine}} + NV_{\text{Mac$$

where  $x_{ij,m} = |x_{i,m} - x_{j,m}|$  distance between electron *i* located at  $x_{i,m}$ and electron j located at  $x_{j,m}$  in wire m, Ewald interaction  $V(x_{ij,m})$  and Madelung energy  $V_{
m Mad}$ 

#### Method

- variational Monte Carlo for Slater-Jastrow trial wave function  $\Psi_T = D[\phi^{\uparrow}(x(\mathbf{R}))]D[\phi^{\downarrow}(x(\mathbf{R}))]e^{J(\mathbf{R})},$
- with orbital spin-up electron  $\phi^{\uparrow}$ , Slater determinant D
- backflow transformation: oordinates of electrons in Slater determinants are replaced by "quasi-particle coordinates" related to actual electron positions by backflow functions consisting of polynomial expansions in electron in-wire separation up to 8th order
- CASINO's Jastrow factors N. D. Drummond, M. D. Towler, and R. J. Needs, Phys. Rev. B 70, 235119 (2004)
- non-reweighted variance minimization N. D. Drummond and R. J. Needs, Phys. Rev. B 72, 085124 (2005) to optimize free parameters of trial wave function,  $5 \times 10^6$  statistically independent steps and 1024 configurations • comparison of variational MC and diffusive MC yields better than 99.9% agreement in ground state energy but less computationally costly

# Correlational energy

interaction energy per electron  $\Delta E = E_{
m g}^2 - E_{
m g}^1$  as difference between biwure and single wire



Symbols represents data for single wire • for  $r_s = 5$  a small peak develops in  $g_{11}$  at  $r = r_s$  when inter-wire distance is reduced to 0.6 a.u. rising with reduction in d

• at distance d = 0.4 a.u.  $g_{11}$  oscillates with period of  $r_s$  rather than with  $r = 2r_s$ 

• similar to  $g_{11}$ , also  $g_{12}$  oscillates at period  $r = r_s$  for  $d \leq 0.4$  a.u. • at  $r_s = 5$  for close proximity of two wires, we interpret change in period as tendency towards formation of Wigner crystal phase

• at  $r_s = 10$  oscillations in both inter- and intra-wire PCFs are enhanced further, inter-wire correlations are comparatively stronger than intra-wire correlations as range of d is significantly smaller than  $r_s$ 

• two kinds of osculations, period of  $r_s$  enveloped by second kind of oscillation due to interplay between intra- and inter-wire correlations

#### Static structure factor



### Summary

- parallel infinitely-thin electron-electron quantum biwire system studied by variational Monte Carlo method
- interaction energy falls off as  $d^{-2}$  for high densities where d is comparable to  $r_s$ , for low densities  $d \ll r_s$  smaller exponent
- pair correlation function shows oscillatory behavior with two frequencies • as wires approach, inter-wire correlations increase while intra-wire correlations decrease
- static structure factor shows a peak at  $2k_F$  at higher densities, a second peak starts to appear at  $4k_F$  when  $r_s = 2$  and d = 0.2 a.u., which is not found for d > 0.2
- for lower densities, first peak completely disappears and the height of the second peak keeps increasing with  $r_s$  and d
- electron-electron biwire system goes into a quasi-Wigner crystalline state at densities higher compared to the case of a single wire
- momentum distribution reveals a Tomonaga-Luttinger (TL) liquid behavior with power law nature near  $k_F$  even in presence of an extra interwire interaction

books and paper: http://www.k-morawetz.de

1. Eur. Phys. J. B 91 (2018) 29, Dependence of structure factor and



ting, works for  $1 < r_s < 20$  too • interaction energy  $\Delta E(d)$  and correlation energy of single wire  $E_c(d) = E_c^1 + C_c^2$  $\Delta E(d)$ , therefore, dependence of correlation energy on wire separation d is similar to  $\Delta E(d)$ 

#### Pair-correlation function

intra-wire (parallel-spin) m = 1, inter-wire (anti-parallel-spin) m = 2





• oscillations in  $g_{12}$  increases, in  $g_{11}$  decreases, as d decreases • correlations between electrons of different wires are build on and intrawire correlations are suppressed as two wires approach

• inter-wire SSF  $S_{12}(k)$  negative for •  $S_{12}(k)$  gets positive before  $4k_F$  and a second peak builds up at  $4k_F$  for d = 0.4 whose height increases as d

#### Sum over spin pairs





open circles: isolated single wire

• for high densities small peak at  $2k_F$  whose height decreases as d becomes smaller (a), (b)

• lower density, correlation effects become more important (c), for  $r_s = 2$ 

- correlation energy on the width of electron wires, Vinod Ashokan, Renu Bala, Klaus Morawetz, and Karem N. Pathak
- 2. Phys. Rev. B 97 (2018) 155147, Conditions where RPA becomes exact in the high-density limit, Klaus Morawetz, Vinod Ashokan, Renu Bala, and Karem N. Pathak
- 3. Phys. Rev. B 101 (2020) 075130, Exact ground-state properties of the one-dimensional electron gas at high density ,Vinod Ashokan, Renu Bala, Klaus Morawetz, and K. N. Pathak
- 4. Phys. Rev. B 104 (2021) 035149, Ground-state properties of electronelectron biwire systems, R. O. Sharma, N. D. Drummond, V. Ashokan, K. N. Pathak, K. Morawetz
- 5. Phys. Rev. B 105 (2022) 115140, Electron correlation and confinement effects in quasi-one-dimensional quantum wires at high density, A. Girdhar, V. Ashokan, N. D. Drummond, K. Morawetz, K. N. Pathak



- first peak in  $g_{11}$  and  $g_{12}$  near  $r = 2r_s$  and  $r = r_s$ , both  $g_{11}$  and  $g_{12}$ oscillates with a period  $2r_s$
- as d reduces, first peak of  $g_{12}$  rises and shifts towards origin while for  $g_{11}$ it shrinks and shifts away from origin
- $g_{12}(r)$  at r = 0 shifts towards zero as d reduces, because with decreasing
- d, electrons in different wires repel each other and  $g_{12}(0)$  becomes smaller

a second peak appears at  $4k_F$  when d is reduced to 0.2 a.u., no such peak in single isolated wire

- increasing  $r_s$ , the  $4k_F$  peak rises while no  $2k_F$  peak  $d \ 1...0.1$
- at  $r_s = 20$  the peak at  $2k_F$  reappears (f)
- peak at  $4k_F$  signals ordered structure, Wigner crystallization at much
- lower densities as compared to the single wire